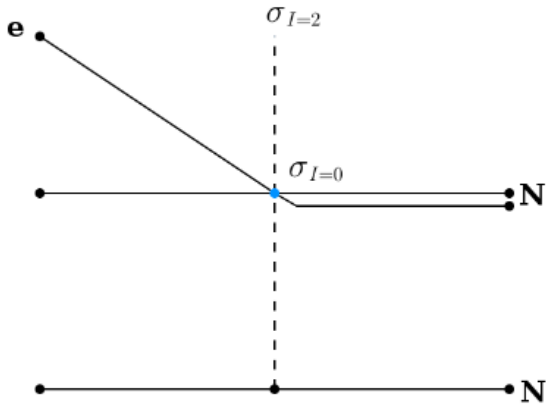
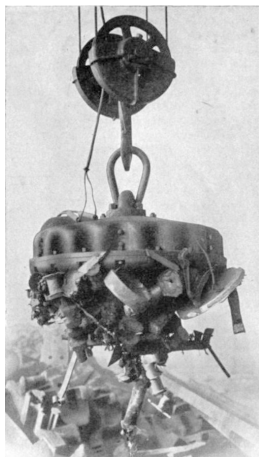
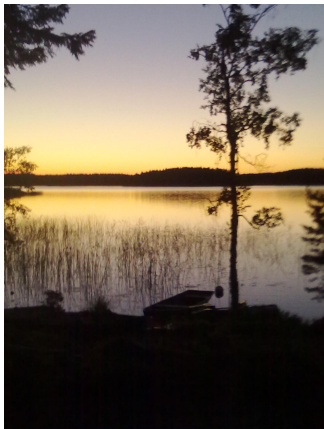


# Working with theory about the Rossi effect



# Invitation



- ▶ No strong magnets are found in nature. (on the surface of earth)
- ▶ Control current to enhance special strong directed magnetism in metals not found in nature.
- ▶ Control electrons to enhance special strong directed nucleon-nucleon force not found in nature.

## *“Low Energy Nuclear Reactions”*

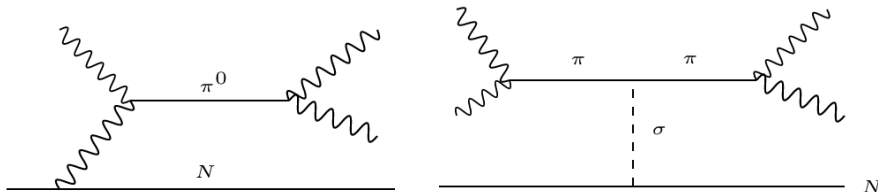
Problems to form a theory around:

- ▶ No strong radiation. If no detection of strong radiation is found together with evidence of isotopic shifts ie the strong force, then everything that creates strong radiation must be forbidden and a theory is formed around what is left.
- ▶ The limited range of the strong force.
- ▶ Nuclides don't come close enough at room temperature to affect each other with the strong force.

Solution presented here:

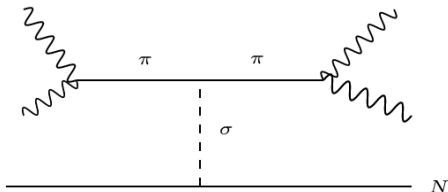
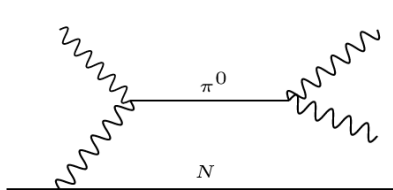
A new special potential of the strong force that is not found (common) in nature.

- ▶ Important feature is electron-nucleon interaction mediated by  $\sigma$  mesons.
- ▶ Releases energy continuously by slowly accelerate nuclides giving them kinetic energy.



# LENR

- ▶ The special potential is a strong force potential triggered by electrons. Hence it does not require long range nucleon-nucleon interaction as a start point.
- ▶ The special conditions that is not natural is that the electrons has to stay near( $10^{-15}$  m scale) the nuclide for a long time while relative spin, velocity and space relation has to be comparable with binding condition of nucleon nucleon interaction.



# Outline

- ▶ Main theory in 3 steps
- ▶ Short on other theories
- ▶ Experiment
- ▶ Comparison theory to experiment
- ▶ Future

# Outline

The main theory is developed in three steps:

1. Develop a strong force potential for nucleon-nucleon(N-N) interaction with the no  $\gamma$  radiation as a requirement.
  2. Enhance this potential by electron- $\sigma$  meson interaction using isospin splitting of the  $\sigma$  meson in nucleons.
  3. Use nucleon polarizability theory to establish the electron nucleon interaction properties based on electromagnetic field component in the interaction.
- Note: Our paper<sup>1</sup> has theory presented in opposite order.

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<sup>1</sup><https://arxiv.org/abs/1703.05249>

## No $\gamma$ problem for fusion

“cold fusion” is a bad word why? Because  $\gamma$  radiation needed for fusion reactions.

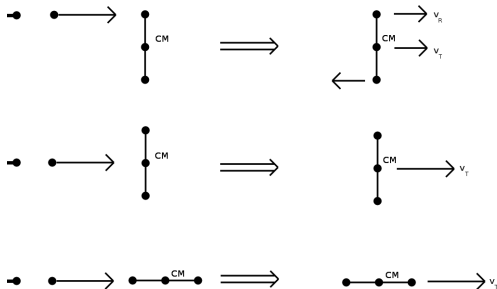
- ▶ Momentum and energy conservation can't be conserved in a pure  $2 \rightarrow 1$  interaction (unless the sum of the initial momentum is 0).
- ▶ Normal fusion has an extra photon to carry away the extra momentum.
- ▶ Nucleon transfer reaction is a  $2 \rightarrow 2$  body reaction. Energy and momentum is allowed to be conserved without extra particles.



## No excited state for nucleon transmission problem

Nucleon transmission reactions solves momentum conservation problem.

However just add a nucleon with a momentum transfer on a nuclide might create an oscillation motion unless the momentum transfer is applied on the center of mass.



## No excited state for nucleon transmission problem

- ▶ Oscillation motion of nucleon inside nuclides=Excited states.
- ▶ Excited states in nuclides de-excite emitting strong radiation (most  $\gamma$ ).
- ▶ Momentum transition to center of mass required for non excited state.

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### Potential from requirement:

- ▶ Nucleon transfer reaction needed.
- ▶ Momentum transfer must be applied on center of mass on both nuclides.
- ▶ Attraction to Mass=Scalar term needed in attractive potential.

## Strong force

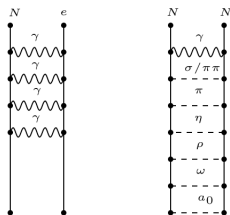
Theoretical need for LENR:

- ▶ Nucleon-nucleon interaction.
- ▶ Nucleon electromagnetic interaction i.e. nucleon-electron/photon interaction.

Energy range	fermion-boson (stable ( $\sim$ point-)particle -interaction particle)	Type	Main experimental need
$\sim 1\text{GeV} >$	quark-gluon	Complete?	decay of hadrons/mesons
keV-MeV	baryon(nucleon)-meson	Effective	Nucleons in nuclides

- ▶ Mesons: 2 quark state separated by spin, charge, parity and quark generation(isospin: $l, \tau$  for generation 1).
- ▶ Baryons: 3 quark state Example:proton  $p(I = 1/2)$ , neutron  $n(I = 1/2)$  and Delta $\Delta(I = 3/2)$
- ▶ Effective theories uses LEC's(low energy constant) to describe phenomena in a lower energy range. Many theories exist depending on choice of approximative formula and problem.

# Strong force



Nucleon nucleon(N-N) interaction:

- ▶ Using complete quark-gluon theory has problem with fermion doubling problem. Also quark theory can't explain the absent of electric dipole moment of the proton and neutron.
- ▶ Nucleon nucleon interaction: Derived from meson exchange between nucleon. This is usually done as a trial function most succesfull theories fits constanst directly to  $r$  space operators.
- ▶ Both 2 and 3 nucleon interaction needed to explain all observations.

## Strong force

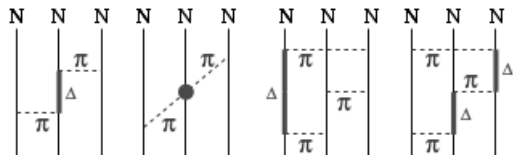
$$V_{L,S,I}(r) = \int \frac{d^3q}{(2\pi)^3} e^{iqr} \frac{g^2}{-q^2 - m^2} O(L, S, I) = -\frac{g^2}{4\pi} \frac{e^{-mr}}{r} O(L, S, I)$$

Meson	mass(MeV/c <sup>2</sup> )	$I (J^P)$	role in N-N potential
$\pi$	138	$1 (0^-)$	Classic long range
$\sigma$	550	$0 (0^+)$	Binding(central+ $L \cdot S$ )
$\omega$	782	$0 (1^-)$	repulsive
$\rho$	770	$1 (1^-)$	$L \cdot S$ interaction
$a_0$	980	$1 (0^+)$	short range
$\eta$	548	$0 (0^-)$	binding

# Strong force

3 nucleon interaction:

- ▶ Internal structure change in the nucleon leads to different potential.
- ▶ Also the exchange particles interact in the space between two nucleons.



## $\sigma$ meson

- ▶ Scalar meson, scalar term=center of mass i.e. property of nucleon potential derived from no  $\gamma$  requirement.
- ▶ In effective field theory:  $\pi\pi$  s-wave resonance.
- ▶ Not natural long range. In one boson exchange potential(NN interaction usually change this term with one include a  $\Delta$ ):

$$V_{NN}^{(\sigma)}(r) = \int \frac{d^3q}{(2\pi)^3} e^{iqr} \frac{g_{\sigma NN}^2}{-q^2 - m_\sigma^2} = -\frac{g_{\sigma NN}^2}{4\pi} \frac{e^{-m_\sigma r}}{r}$$

with  $m_\sigma \simeq 550$  MeV. Compare to EM potential from photon:

$$\frac{q}{4\pi} \frac{e^{-m_\gamma r}}{r}$$

with  $m_\gamma = 0$ .



## $\sigma$ meson isospin

$\sigma$  interaction properties: The  $\sigma$  meson is a phase shift in  $\pi\pi$  scattering separated by isospin states.

- ▶ From  $\pi\pi$  scattering<sup>2</sup>:  $m_{\sigma_{I=0}} = 36.77m_{\pi}^2$  and  $m_{\sigma_{I=2}} = -21.62m_{\pi}^2$
- ▶ Hadron interaction properties: Proton and neutron isospin half state  $\rightarrow m_{\sigma}$  in OBE potential is mixed:

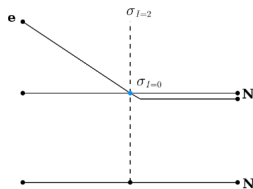
$$m_{\sigma} = \sqrt{m_{\sigma_{I=0}}^2 - m_{\sigma_{I=2}}^2} = 543 \cong 550$$

- ▶ Idea for step 2: electron-nucleon interaction enhance the range of the  $\sigma$  part of a N-N potential.

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<sup>2</sup>G. Colangelo, J. Gasser, H. Leutwyler,  $\pi\pi$  scattering, arxiv:hep-ph/0103088v1

## $\sigma$ meson isospin



- ▶ Why? Electron is isospin 0 state interaction with  $\sigma_{I=2}$  is suppressed compared to  $\sigma_{I=0}$  which would increase the range of the  $\sigma$  part of the N-N potential.
- ▶  $\sigma$ -electron interaction needs electron in nucleon since the interaction range of  $\sigma_{I=0}$  is short.
- ▶ The  $\sigma$ -electron interaction doesn't equal a binding (by theory) therefore the electron must be kept in place by force.
- ▶ The electron must be able to spinflip (hyperfine structure) to extract  $\sigma_{I=0}$  energy out of the nucleon while stay in place.
- ▶ Hyperfine structure interaction is in the energy range of meV while the potential need to absorb  $\sigma$  carried energy in the MeV range. This implies a long startup time.

## Needed electron nucleon interaction

- ▶ Needed  $\sigma$ -electron interaction not complete yet by relative spin, velocity and space relation.
- ▶ If  $\sigma$ -electron interaction can't form bound states, the distance between electron and nucleon must be in the  $10^{-15}$  m(fm) range.
- ▶ At a atomic level this means special condition in which the electron “stay” inside the nucleon.
- ▶ Electron nucleon interaction that would change the internal structure of the nucleon needed. So that the nucleon would correspond to a new bound state in another nuclide.

# Nucleon polarizability

- ▶ What: Electromagnetic(EM) interaction of nucleon besides basic coulomb and magnet interaction i.e. internal structure changes.
- ▶ Compare to the 3 nucleon force where the internal structure changes affect the potential derived from the nucleon.
- ▶ Goal: Find EM interaction of e-N system that corresponds to binding condition of N-N force.
- ▶ Why? This would set the included nuclides in the new ground state binding condition after a nucleon transfer.
- ▶ Binding condition for particle systems are calculated by adding a negative term to the kinetic hamiltonian.
- ▶ Theoretical work have problems with choice of type of effective field approximation.

# Nucleon polarizability

- ▶ Example coulomb interaction hamiltonian for atomic physics ( $V < 0$  electromagnetic binding):

$$H\Psi = E\Psi \rightarrow (-k\nabla^2 + V)\Psi = E\Psi$$

- ▶ Polarizability calculated in perturbation theory by adding effective hamiltonians that are divided according to spacetime derivatives<sup>(i)</sup> of the EM field

$$H = E_0 - \sum H_{\text{eff}}^{(i)}$$

- ▶ We look after conditions  $H_{\text{eff}} > 0$ .

# Nucleon polarizability

Advance equations from perturbation theory<sup>3</sup>:

$$\begin{aligned} H_{\text{eff}}^{(2)} &= -\frac{1}{2}4\pi (\alpha_{E1}\bar{E}^2 + \beta_{M1}\bar{H}^2) \\ H_{\text{eff}}^{(3)} &= -\frac{1}{2}4\pi [\gamma_{E1E1}\bar{\sigma} \cdot (\bar{E} \times \dot{\bar{E}}) + \gamma_{M1M1}\bar{\sigma} \cdot (\bar{H} \times \dot{\bar{H}}) \\ &\quad - 2\gamma_{M1E2}E_{ij}\sigma_i H_j + 2\gamma_{E1M2}H_{ij}\sigma_i E_j] \\ H_{\text{eff}}^{(4)} &= -\frac{1}{2}4\pi (\alpha_{E1\nu}\dot{\bar{E}}^2 + \beta_{M1\nu}\dot{\bar{H}}^2) - \frac{1}{12}4\pi (\alpha_{E2}E_{ij}^2 + \beta_{M2}H_{ij}^2) \end{aligned} \quad (1)$$

$\alpha_x, \beta_x, \gamma_x$  = polarizability constants.

$\sigma$  = Pauli spin matrices of the nucleon

E and H are components of the electromagnetic fields.

$$E_{ij} = \frac{1}{2}(\nabla_i E_j + \nabla_j E_i) \text{ (same for } H_{ij}\text{)}$$

Note: The third order perturbation is called spin polarizability and is not included in an classic static EM field.

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<sup>3</sup>F.Hagelstein, R.Miskimen and V.Pascalutsa, "Nucleon Polarizabilities: from Compton Scattering to Hydrogen Atom," Prog. Part. Nucl. Phys. bf 88 (2016) 29 [arXiv:1512.03765 [nucl-th]].

## Values

Theoretical and experimental values of the proton and neutron static dipole, quadrupole and dispersive polarizabilities. The units are  $10^{-4} fm^3$  (dipole) and  $10^{-4} fm^5$  quadrupole.

	$\alpha_{E1}$	$\beta_{M1}$	$\alpha_{E2}$	$\beta_{M2}$
Proton				
B $\chi$ PT Theory <sup>4</sup>	$11.2 \pm 0.7$	$3.9 \pm 0.7$	$17.3 \pm 3.9$	$-15.5 \pm 3.5$
Experiment(PDG <sup>5</sup> )	$11.2 \pm 0.4$	$2.5 \pm 0.4$		
Neutron				
B $\chi$ PT Theory	$13.7 \pm 3.1$	$4.6 \pm 2.7$	$16.2 \pm 3.7$	$-15.8 \pm 3.6$
Experiment(PDG)	$11.8 \pm 1.1$	$3.7 \pm 1.2$		

<sup>4</sup>V.Lensky and V.Pascalutsa, "Predictive powers of chiral perturbation theory in Compton scattering off protons," Eur. Phys. J. C 65 (2010) 195 [arXiv:0907.0451 [hep-ph]].

<sup>5</sup>C. Patrignani et al.(Particle Data Group), Chin. Phys. C, 40, 100001 (2016)

# Values

Theoretical values of the proton and neutron static dispersive polarizabilities. The units are  $10^{-4} \text{ fm}^5$ .

	$\alpha_{E1\nu}$	$\beta_{M1\nu}$
Proton		
B $\chi$ PT Theory	$-1.3 \pm 1.0$	$7.1 \pm 2.5$
Neutron		
B $\chi$ PT Theory	$0.1 \pm 1.0$	$7.2 \pm 2.5$



## Values

Theoretical and experimental values of the proton and neutron static spin polarizabilities. The units are  $10^{-4} \text{ fm}^4$ .

	$\gamma_{E1E1}$	$\gamma_{M1M1}$	$\gamma_{E1M2}$	$\gamma_{M1E2}$
Proton				
B $\chi$ PT Theory	$-3.3 \pm 0.8$	$2.9 \pm 1.5$	$0.2 \pm 0.2$	$1.1 \pm 0.3$
MAMI 2015 <sup>6</sup>	$-3.5 \pm 1.2$	$3.16 \pm 0.85$	$-0.7 \pm 1.2$	$1.99 \pm 0.29$
Neutron				
B $\chi$ PT Theory	$-4.7 \pm 1.1$	$2.9 \pm 1.5$	$0.2 \pm 0.2$	$1.6 \pm 0.4$

- ▶ Sign of  $\gamma_{E1M2}$  visualize problem with different effective field theories:

$O(p^4)_b$	$O(\epsilon^3)$	$O(p^4)_a$	K-Matrix	HDPV
0.7	1.0	0.2	-1.8	-0.02
DR	$L_\chi$	HB $\chi$ PT	B $\chi$ PT	MAMI 2015
-0.02	-0.7	$-0.4 \pm 0.4$	$-0.2 \pm 0.2$	$-0.7 \pm 1.2$

<sup>6</sup>P.P.Martel et al. [A2 Collaboration], "Measurements of Double-Polarized Compton Scattering Asymmetries and Extraction of the Proton Spin Polarizabilities," Phys. Rev. Lett. 114 (2015) [arXiv:1408.1576 [nucl-ex]] ▶

## Polarizability binding conditions

- ▶ The  $H_{eff} < 0$  condition has to be valid for the full equation, so that there can't be an extra E field if the B field condition is fulfilled plus the opposite.
- ▶ Define variable  $x_{L,T} = \dot{\vec{E}}_{L,T} / \bar{E} \cdot \hat{\vec{E}} / \dot{\vec{E}}$  to get two differential equations:

$$\alpha_{E1} \pm \gamma_{E1E1} x_T + \alpha_{E1\nu} x_T^2 \quad (2)$$

The  $\pm$  sign is determined by the direction between the vectors  $\bar{\sigma}$  and  $(\bar{E} \times \dot{\vec{E}})$ .

$$\alpha_{E1} + \alpha_{E1\nu} x_L^2 \quad (3)$$

## Polarizability binding conditions for electric field

Calculations for theoretical values from  $B\chi PT$  gives the  $H_{eff} > 0$  ranges:

Nucleon	$\text{sgn} \bar{\sigma} \cdot (\bar{\mathbf{E}} \times \dot{\bar{\mathbf{E}}})$	$x = \dot{\bar{\mathbf{E}}}/\bar{\mathbf{E}}$ range (fm)
p	+	$x_T < -2$ $x_T > 4.5$
p	-	$x_T < -4.5$ $x_T > 2$
p	0	$x_L^2 > 0.11$
n	+	$3.1 < x_T < 44$
n	-	$-44 < x_T < -3.1$
n	0	-

## Polarizability binding conditions for magnetic field

- ▶  $\beta_{M2} < 0$  gives  $H_{eff} > 0$  at a center of magnetic quadrupole.

For combined electric and magnetic fields define  $x = \sigma_i E_j / H_{ij}$  and  $E_j$  as  $E \sin \theta$  (with  $\theta$  the angle between dimension  $j$  and the plane defined by  $i$  and  $k$ ). This gives the second order equation:

$$\beta_{M2}/6 + 2 \sin \theta \gamma_{E1M2} x + x^2 \alpha_{E1} \quad (4)$$

$x=0$  always gives  $H_{eff} > 0$  values. The relation to have  $H_{eff}$  values with different sign are given by::

$$\frac{6 \gamma_{E1M2}^2 \sin^2 \theta}{\beta_{M2}} = \alpha_{E1} \quad (5)$$

$$H_{\text{eff}} > 0$$

The three conditions for  $H_{\text{eff}} > 0$ :

- ▶ A center of a magnetic quadrupole which also allows for a weak electric field.
- ▶ Two ranges from the parameter  $\dot{\vec{E}}_{L,T}/\bar{E}$ :
- ▶ The  $x_L^2 > 0.11$  range has  $\dot{\vec{E}}$  in the direction of  $\bar{E}$ .
- ▶ The  $x_T$  ranges has  $\dot{\vec{E}}$  perpendicular to  $\bar{E}$ , this means a circular motion of the electron around the nucleon.

## Atomic states

- ▶ Combine the short range need from the  $\sigma_{I=0}$  mass with the spin polarizabilities yields special conditions on the nucleon electron relation.
- ▶ The energy gets released by accessing the  $\sigma$  meson part of the nucleon by polarizability relations and sent away by classic electromagnetic interaction.
- ▶ Due to the long range that the new potential is suppose to have, the electron has to have a stable position near the nuclide(in the fm range).
- ▶ Only atomic binding to have electron at nuclide is s-state atomic bindings however average distance is in the order of  $10^{10}$  m.
- ▶ Solutions:
  1. Pressure would make the electron come closer to the nucleon.
  2. The nuclides create a positive ion current that follows an electron current.

## Atomic states

- ▶ The  $x_L$  solution can't be combined with the long time requirement since the interaction is a linear motion.
  - ▶ The magnetic quadrupole solution is compatible with both solution 1 and 2.
  - ▶ The  $x_T$  solution could be combined with solution 1 and 2 in special conditions:
1. Non s-state atomic binding would need extreme pressure to fulfill the range condition. For a s-state element the electron nuclide relation is approximately a rotation if the center of mass does not equal the center of charge. However the nucleon spin is in s-state perpendicular to the  $(\vec{E} \times \dot{\vec{E}})$  vector not aligned. The solution is to have an electron that transforms between a state with aligned nucleon spin and near nucleus condition. The right conditions are found for s- $d_{z^2}$  overlaps.
  2. For nucleon current that follows electron current the nucleon spin must precess around the electron.

## Other theories

New particle theories:

- ▶ Dark matter theories.
- ▶ Low energy virtual particles theory. I and Rossi are also making the hypothesis of the possibility that the temperatures of the plasma can reach the mass of , waves in fields that could annihilate without emitting high energy radiations because of the low energy.

Problem: Strong evidence of isotopic shifts requires link to the strong force at some point.

- ▶ Multi particle binding would explain the no  $\gamma$  condition with binding a lot of particles to the nucleon instead of one.

Problem: Electromagnet interaction strongly enhance one photon couplings.



# Experiment

Observations<sup>789</sup>:

- ▶ Energy production without strong radiation.
- ▶ Isotopic shifts
- ▶ Positive ion current

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<sup>7</sup>[http://www.elforsk.se/Global/Omv%C3%A4rld\\_system/filer/LuganoReportSubmi](http://www.elforsk.se/Global/Omv%C3%A4rld_system/filer/LuganoReportSubmi)

<sup>8</sup>K. A. Alabin, S. N. Andreev, A. G. Parkhomov. Results of Analyses of the Isotopic and Elemental Composition of Nickel- Hydrogen Fuel Reactors.

<https://drive.google.com/file/d/0B5Pc25a4cOM2cHBha0RLbUo5ZVU/view>

<sup>9</sup><https://arxiv.org/abs/1703.05249>

## No $\gamma$ radiation

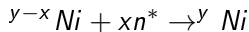
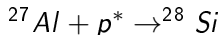
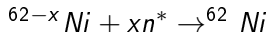
- ▶ Observed: Gamma radiation less than background level.
- ▶ For a detector  $\sim 0.5$  m away this means  $10^4 - 10^6$   $\gamma$ /s for  $\gamma$  energies  $\sim 100$  keV to  $\sim 10$  MeV.
- ▶ Observed:  $10^{12} - 10^{15}$  transfer reactions/s. If each reaction creates  $\sim 1$  MeV of energy.
- ▶ Creation of some radioactive nuclides that almost does not produce  $\gamma$  radiation still possible.
- ▶ Example:

${}^6\text{He}$ :  $\sim 10^{12}$  produced/s possible, above this secondary x-ray from  $\beta^-$  radiation should be detectable.

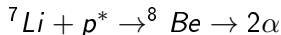
${}^{59}\text{Ni}$ :  $\sim 10^{20}$  produced/s possible. Above this rate 511 keV  $\gamma$  rays from positron annihilation should be at background level.

# Isotopic shifts

Main detected isotopic shifts:



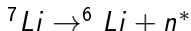
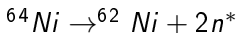
where  $p^*$  and  $n^*$  mean a bound nucleon. Also possible observed is:



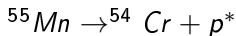
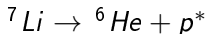
(If  $p$  is a free proton this would create measurable  $\gamma$  radiation above background level but not with bound)

# Isotopic shifts

Neutron sources:



Proton sources:



- ▶ Note that the transmitted nucleon has to be absorbed into a lower energy state i.e. isotopic shifts only happens for energy release reactions.

## Positive ion current

New experimental observation: Li/H ratio in plasma is related to output energy.

Output power is created when negative ions changes to positive ion kinetic energy.

Neutral plasma  $\rightarrow$  number and speed of positive and negative ions that enters the plasma are the same.

COP: Kinetic energy of positive ions/kinetic energy of negative ions.

Non relativistic kinetic energy:

$$\sum \frac{m_+ v_+^2}{2} / \sum \frac{m_- v_-^2}{2}$$

- ▶ Neutral plasma gives:  $\sum v_+^2 = \sum v_-^2$
- ▶ COP is related to  $m_+/m_-$  i.e. in the range  $m_{Li}/m_e = 14000$  to  $m_H/m_e = 2000$ .
- ▶ Measured COP in the doral test are in the range of thousands. Li/H ratio are reduced with the COP.

## Important atomic states

Experimental observation of needed elements is in agreement with the theoretical requirement of atomic states.

- ▶ Free s-state electron elements needed to have spin flip electrons in nucleon.
- ▶ Free  $d_{z^2}$  electron elements needed to have nucleon spin perpendicular to electron.
- ▶  $d_{z^2}$ -s overlap needed

$d_{z^2}$  electron elements:

- ▶ Nickel group i.e. Nickel, Palladium, Platinum.

Free s-state electron elements:

- ▶ Hydrogen
- ▶ Alkalimetals: Lithium, Sodium, Potassium, Cesium.
- ▶ Some other metals: nickel, platinum, niobium, molybdenum, ruthenium, rhodium, and chromium.

## Experiment-theory comparison summation

- ▶ No strong radiation  $\rightarrow$  continuous kinetic energy release of nuclides.
- ▶ (No  $\gamma$ ) Momentum transfer is applied on center of mass  $\rightarrow$  potential is mediated by scalar meson ( $\sigma$ )
- ▶ Long range  $\sigma$  potential not natural. Created by special electron-nucleon interaction.
- ▶ One electron in s-state elements are needed: spin flip electron in the nucleon needed.
- ▶ One free electron in  $d_{z^2}$  shell elements needed: Tilt electron in right position compares to binding condition of polarizability.
- ▶ Plasma between Ni rods: Nickel creates  $\sigma_{l=2}$  potential that drags protons and Lithium ions through air.

# Future

Experimental to do:

- ▶ Important atomic states must be examined better. For example by doing isotopic shift measurement in slices.
- ▶ Measure a start time for the reaction to compare the meV spinflip interaction to the MeV nucleon transfer reaction.
- ▶ Find evidence for more possible isotopic shifts.
- ▶ Take  $\alpha/\beta$  radiation spectrum to fit to theory.
- ▶ Exact numbers of output power compared to H/Li ratio in plasma.
- ▶ Hydrogen in plasma. Free hydrogen or proton in a long range nucleon transfer reaction?(would be possible to measure by the mass of the proton)

Not LENR experiment:

- ▶ Better polarizability measurement to confirm the theoretical values.



# Future

Theoretical to do:

- ▶ More exact theory for electron- $\sigma$  interaction needed.
- ▶ Theory for multinucleus transfer reactions. Especially  $\alpha$  clusters.(Deuteron-Palladium systems)
- ▶ Detailed study of available atomic states.
- ▶ A unified theory for the effective theory range of the strong force. Fermion doubling problem needs to be solved for this(I know a solution but it isa long proof).